

Tumbling in an extensional flow

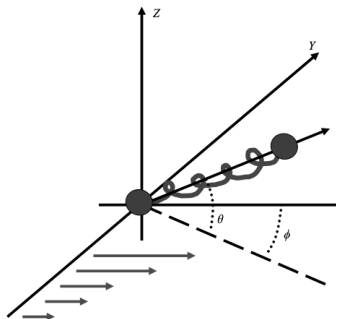
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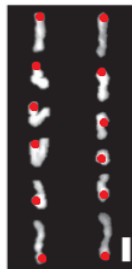
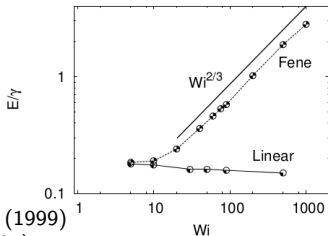
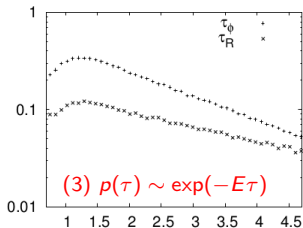
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Polymer dynamics in shear flow



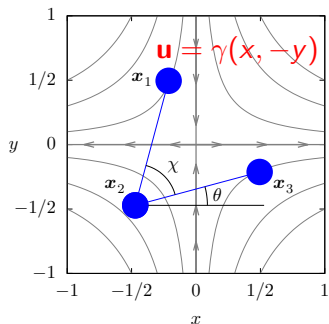
(2) shear flow: $\mathbf{u} = \gamma(0, y, 0)$



(4) DNA tumbling in shear flow
 $L \approx 7\mu\text{m}$

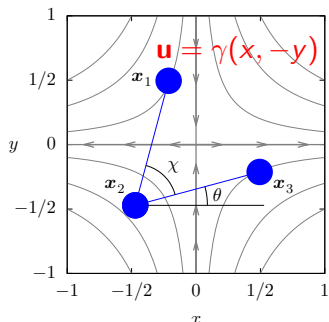
- (1) Smith, Babcock and Chu. Science. (1999)
- (2) Chertkov, et al. J. Fluid Mech. (2005)
- (3) Celani, Puliafito and Turitsyn. Europhys. Lett. (2005)
- (4) Gerashchenko and Steinberg. Phys. Rev. Lett. (2006)

The trumbbell model



- ① rigid, inertialess beads and rods
- ② frame of reference at \mathbf{x}_c : $\mathbf{r}_\nu = \mathbf{x}_c - \mathbf{x}_\nu, \nu = 1, 2, 3$
- ③ coordinates $\mathbf{q} = \{\theta, \chi\}$ (in 2D)
- ④ neglect hydrodynamical interactions
- ⑤ dilute solution

The trumbbell model



Forces at play:

- Stokes' drag with coefficient ζ
- rigidity constraint (fixed length ℓ)
- restoring potential $\phi = \mu(\chi - \pi)^2/2$
- thermal fluctuations

Langevin equations:

$$m\ddot{\mathbf{x}}_\nu = -\zeta[\dot{\mathbf{x}}_\nu - \mathbf{u}(\mathbf{x}_\nu, t)] + \mathbf{t}_\nu + \mathbf{f}_\nu + \sqrt{D}\boldsymbol{\eta}_\nu(t)$$

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Diffusion Equation for p.d.f. $\psi(\theta, \chi; t)$

P.d.f. ψ satisfies (Hassager, JCP (1974)):

$$\frac{\partial \psi}{\partial t} = -\frac{\partial}{\partial \mathbf{q}^i} \left\{ G^{ij} \left[\left(\kappa^{kl}(t) \mathbf{r}'_{\nu} \frac{\partial r_{\nu}^k}{\partial \mathbf{q}^j} - \frac{1}{\zeta} \frac{\partial \phi}{\partial \mathbf{q}^j} \right) \psi - \frac{KT}{\zeta} \sqrt{h} \frac{\partial}{\partial \mathbf{q}^j} \left(\frac{\psi}{\sqrt{h}} \right) \right] \right\}$$

where

$$\kappa^{ij}(t) = \partial^j u^i(\mathbf{x}_c(t), t), \quad H^{ij} = \sum_{\nu, k} \frac{\partial r_{\nu}^k}{\partial \mathbf{q}^i} \frac{\partial r_{\nu}^k}{\partial \mathbf{q}^j}, \quad G = H^{-1}, \quad h = \det H$$

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In 2D:

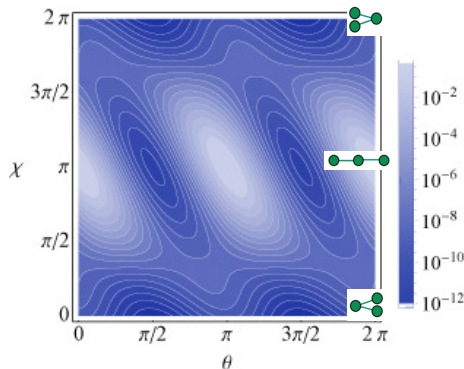
$$\begin{aligned} \dot{\theta}(t) &= V_{\theta} + \sqrt{D_{\theta\theta}} \xi_{\theta}(t), \\ \dot{\chi}(t) &= V_{\chi} + \frac{D_{\theta\chi}}{\sqrt{D_{\theta\theta}}} \xi_{\theta}(t) + \sqrt{D_{\chi\chi} - \frac{D_{\theta\chi}^2}{D_{\theta\theta}}} \xi_{\chi}(t) \end{aligned} \quad (1)$$

where $\xi_{\theta}(t), \xi_{\chi}(t)$ are independent white noises.

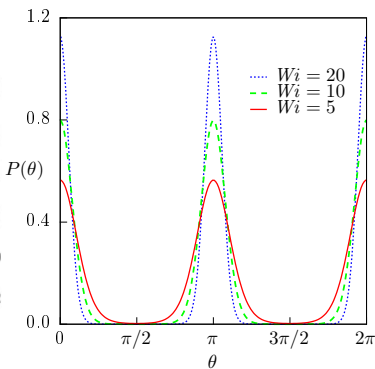
Stationary statistics with $\mathbf{u} = \gamma(x, -y)$

$$\psi_{\text{st}} = N \sqrt{4 - \cos^2 \chi} \exp \left[\frac{Z Wi (2 \cos \chi - 1) \cos(2\theta + \chi)}{3} - \frac{Z(\pi - \chi)^2}{2} \right]$$

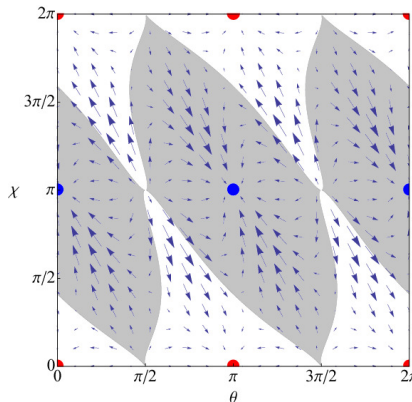
where $Wi = \gamma \zeta \ell^2 / \mu$, $Z = \mu / KT$, $Pe = Z Wi$.



$\ln \psi_{\text{st}}$ for $Wi = 5$, $Z = 1$



marginal p.d.f. of θ

Vector plot ($Wi = Pe = \infty$)

$$P_0 : (\theta = n\pi, \chi = 0, 2\pi)$$

$$\lambda_0^{(1)} = -2\gamma, \lambda_0^{(2)} = -4\gamma/3$$

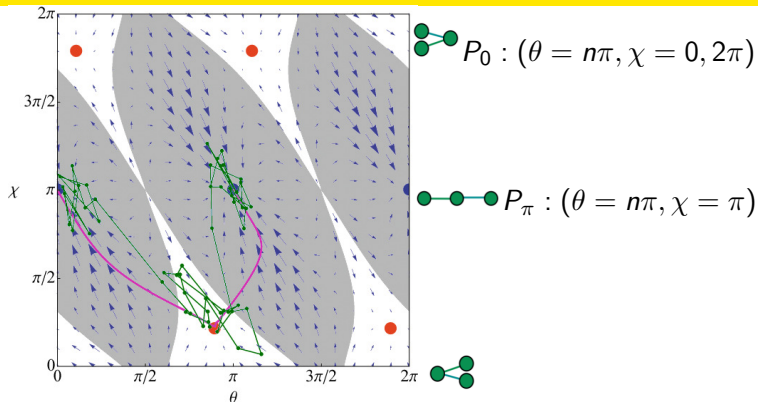
$$P_\pi : (\theta = n\pi, \chi = \pi)$$

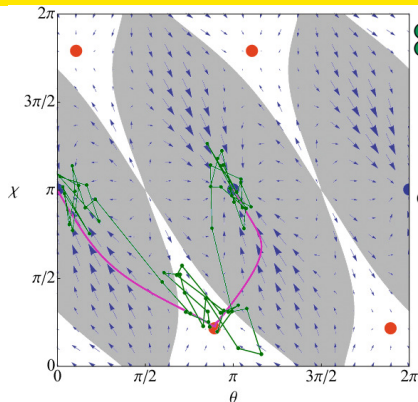
$$\lambda_\pi^{(1)} = -4\gamma, \lambda_\pi^{(2)} = -2\gamma.$$

$$\frac{\lambda_\pi^{(1)}}{\lambda_0^{(1)}} = 2 \text{ for all } \gamma$$

$$\dot{\theta}(t) = -\frac{\gamma}{2(4 - \cos^2 \chi)} [7 \sin(2\theta) + 4 \cos(2\theta + \chi) \sin \chi - \sin(2(\theta + \chi))]$$

$$\dot{\chi}(t) = -\frac{4\gamma}{2 + \cos \chi} [\sin \chi \cos(2\theta + \chi)]$$

Stationary statistics ($Wi < Pe = \infty$)

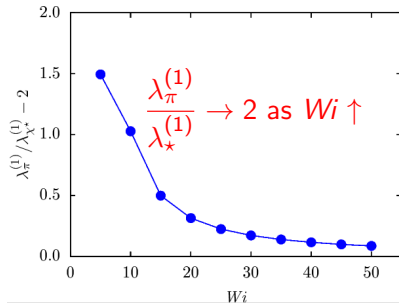
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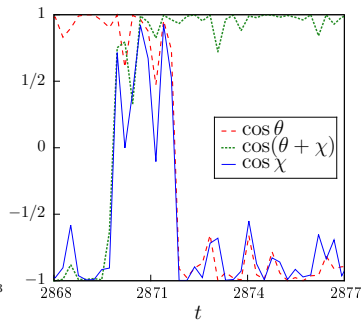
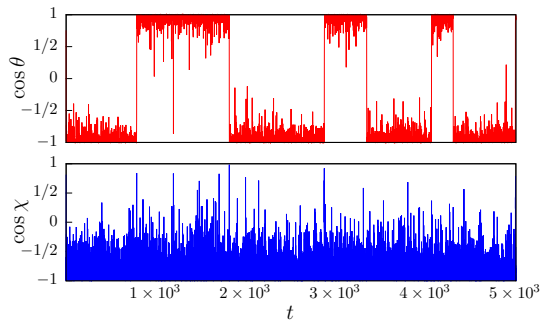
Stable points $P_\star \rightarrow P_0$ as $Wi \rightarrow \infty$

Green curve: sample trajectory
 Pink curve: a minimum energy path
 connecting (π, π) and $(0, \pi)$



Tumbling Dynamics

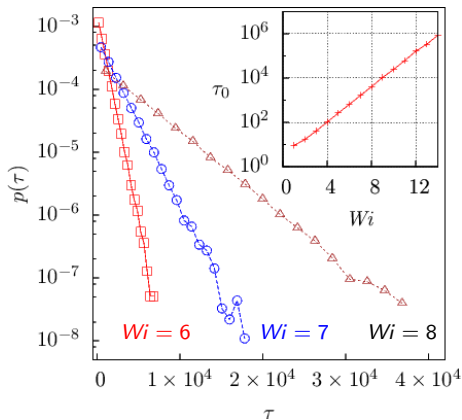
Time series of (1) for $Wi = 6, Z = 1$



Movie: 

P.d.f. of tumbling time τ

$$p(\tau) \propto \exp(-\tau/\tau_0) \text{ for } \tau \gg \tau_0$$

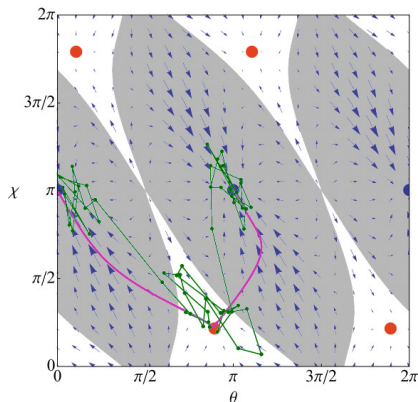


$$\tau_0 \propto \exp(Wi)$$

different from
tumbling time
in shear!

Remarks:

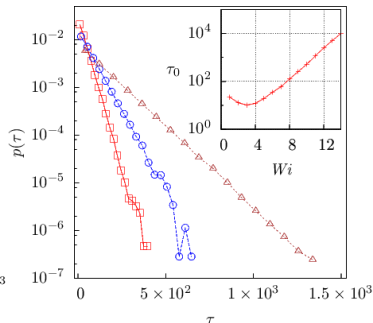
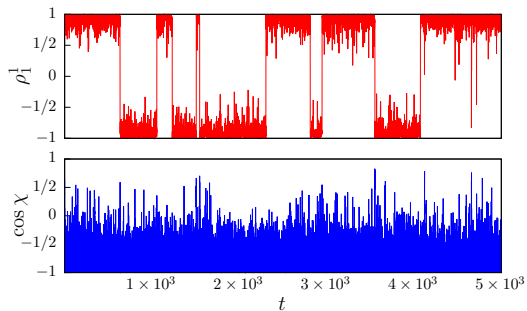
- ① tumbling motion = exit from an attractive domain (large deviations theory)
- ② $p(\text{exit}) = \text{exponential}$ with \uparrow mean as noise \downarrow
- ③ Minimum energy paths (most probable transition path) show tumbling-through-folding motion
- ④ reminiscent of buckling instabilities, but the dynamics are different¹



¹ Lindner and Shelley. In Fluid-Structure Interactions in Low Reynolds-Number Flows. (2014)

Trumbbell model in 3D

Flow: $\mathbf{u} = \gamma(x, -y, 0)$



Coordinates: $\mathbf{q} =$ three Euler angles and $\chi \rightarrow$ singularity in equations

Equivalent formulation: $\dot{\rho}_1 = \dot{\mathbf{x}}_1 - \dot{\mathbf{x}}_2$ and $\dot{\rho}_3 = \dot{\mathbf{x}}_3 - \dot{\mathbf{x}}_2$

\rightarrow system of Stratonovich SDEs

Conclusions

- 1 Semi-flexible polymers perform tumbling-through-folding motion in an extensional flow
- 2 Two dimensional analysis \leftrightarrow planar 3D flow.
- 3 Tumbling time $p(\tau) \propto \exp(-\tau/\tau_0)$ with $\uparrow \tau_0$ as $\uparrow Wi$
- 4 Use of large deviations theory

Questions: Increase degrees of freedom? Effect of dynamics on flow?
Hydrodynamic interactions?

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Questions: Increase degrees of freedom? Effect of dynamics on flow?
Hydrodynamic interactions?

Thank you for your attention!

References

- 1 A. Celani, A. Pulifiato and K. Turitsyn. *Polymers in linear shear flows: A numerical study*. Europhys. Lett., **70**(4), 464-470 (2005).
- 2 M. Chertkov, I. Kolokolov, V. Lebedev and K. Turitsyn. *Polymer statistics in a random flow with mean shear*. J. Fluid Mech. **531**. (2005)
- 3 S. Gerashchenko and V. Steinberg. *Statistics of Tumbling of a Single Polymer in Shear Flow*. Phys. Rev. Lett. **96**, 038304 (2006).
- 4 O. Hassager. *Kinetic theory and rheology of bead-rod models for macromolecular solutions*. J. Chem. Phys. **60**, 5 (1974).
- 5 D. Smith, H. Babcock and S. Chu. *Single polymer dynamics in steady shear flow*. Science. **283** (1999)

Appendix

Drift coefficients of the Fokker–Planck equation in 2D:

$$\begin{aligned}
 V_{\theta} &= -\frac{6KT \sin \chi}{\zeta \ell^2 (2 - \cos \chi)(2 + \cos \chi)^2} - \frac{3\mu(\pi - \chi)}{\zeta \ell^2 (2 + \cos \chi)} - \\
 &\quad \frac{1}{4 - \cos^2 \chi} \{ [2 \sin \chi \cos^2(\theta + \chi) + \cos \theta (4 \sin \theta - \cos \chi \sin(\theta + \chi))] \kappa^{11}(t) + \\
 &\quad [2 \cos \chi \sin^2(\theta + \chi) - \sin \theta (2 + \cos \chi) \sin(\theta + \chi) + 4 \sin^2 \theta] \kappa^{12}(t) + \\
 &\quad [-2 \cos \chi \cos^2(\theta + \chi) + \cos \theta (2 + \cos \chi) \cos(\theta + \chi) - 4 \cos^2 \theta] \kappa^{21}(t) + \\
 &\quad [2 \cos \theta (\sin(\theta + \chi) - 2 \sin \theta) + \cos \chi (\sin \theta \cos(\theta + \chi) - \sin(2(\theta + \chi)))] \kappa^{22}(t) \} \\
 V_{\chi} &= \frac{12KT \sin \chi}{\zeta \ell^2 (2 - \cos \chi)(2 + \cos \chi)^2} + \frac{6\mu(\pi - \chi)}{\zeta \ell^2 (2 + \cos \chi)} + \frac{\sin \chi}{2 + \cos \chi} \{ [1 - 2 \cos(2\theta + \chi)] \kappa^{11}(t) \\
 &\quad - 2 \sin(2\theta + \chi) [\kappa^{12}(t) + \kappa^{21}(t)] + [1 + 2 \cos(2\theta + \chi)] \kappa^{22}(t) \}
 \end{aligned}$$

Noise coefficients:

$$D_{\theta\theta} = \frac{12KT}{\zeta \ell^2 (4 - \cos^2 \chi)}, \quad D_{\theta\chi} = -\frac{6KT}{\zeta \ell^2 (2 + \cos \chi)}, \quad D_{\chi\chi} = \frac{12KT}{\zeta \ell^2 (2 + \cos \chi)}$$